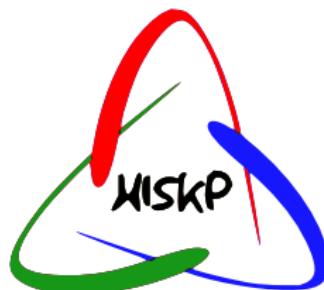
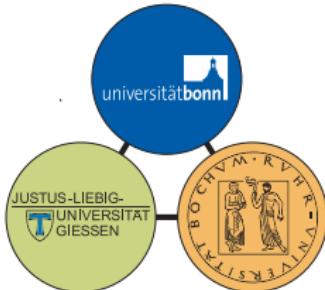


# Measurement of the beam asymmetry $\Sigma$ in $\pi^0$ - and $\eta$ -photoproduction

Farah Noreen Afzal  
for the  
CBELSA/TAPS collaboration

HISKP, University of Bonn

20.02.2015

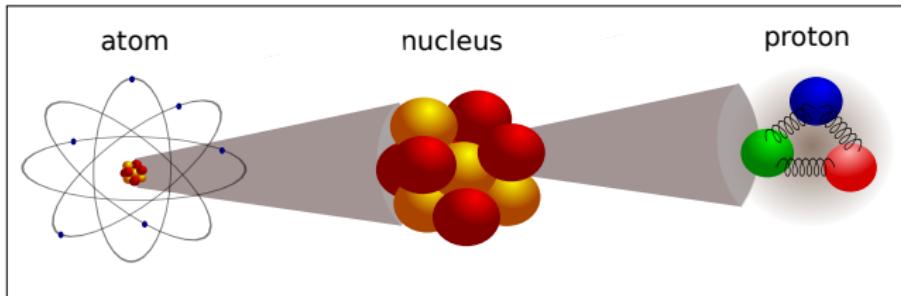


# Outline

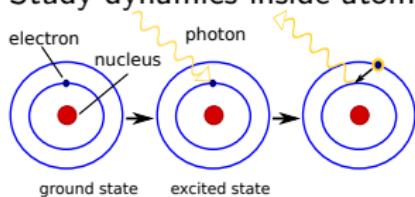
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- 1 Motivation
- 2 The CBELSA/TAPS experimental setup
- 3 Event selection
- 4 Determination of the beam asymmetry  $\Sigma$
- 5 Preliminary results

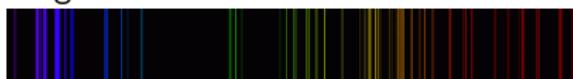
# Why baryon spectroscopy?



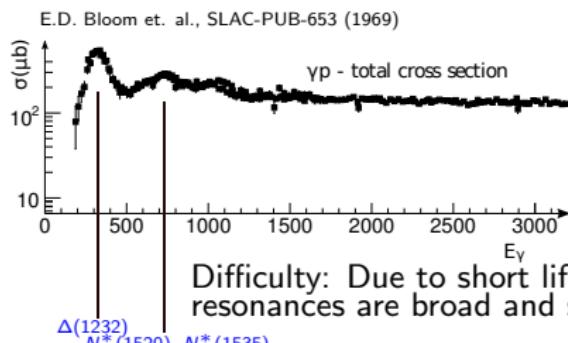
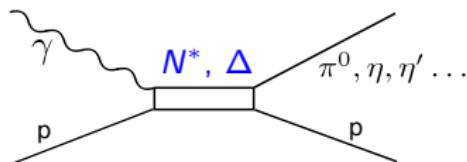
Study dynamics inside atom



Argon:

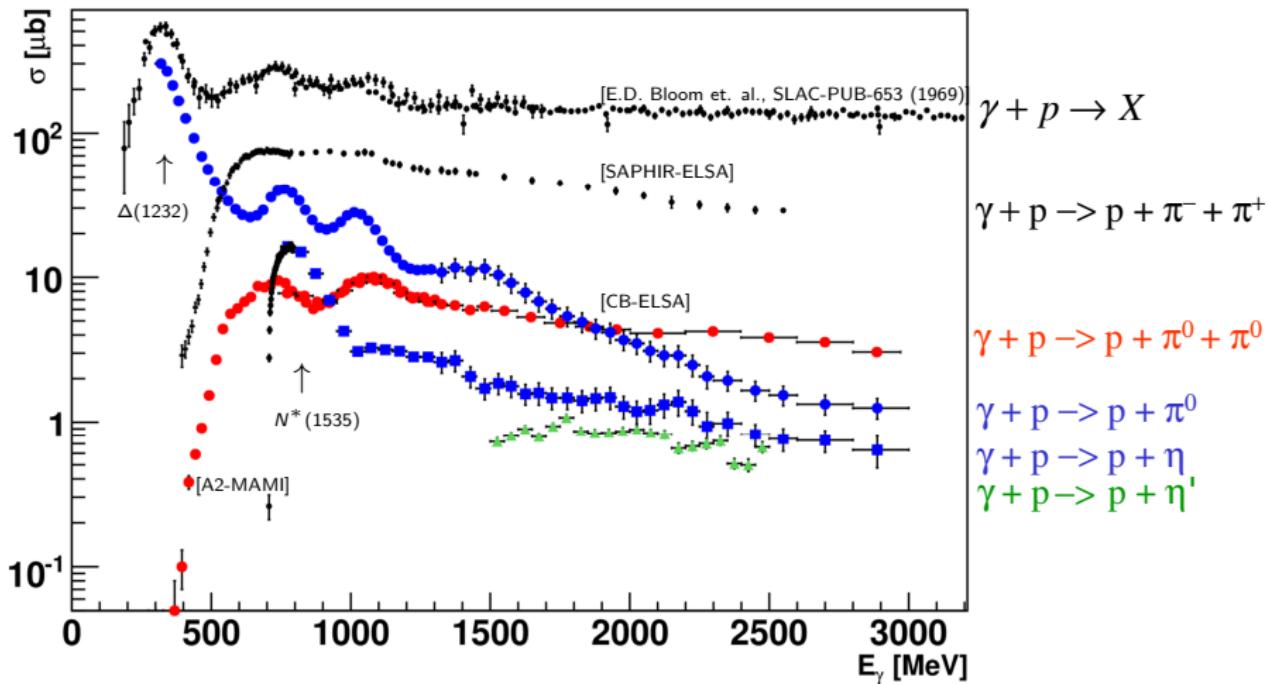


Study dynamics of constituents  
inside the nucleon



# Photoproduction reactions

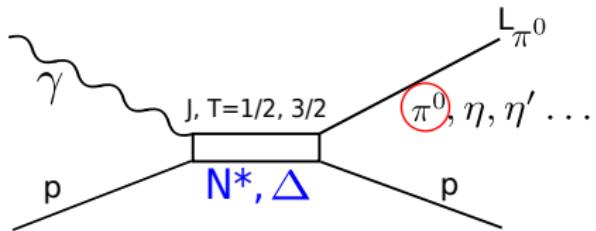
Study of different reaction channels gives access to different resonant structures  
⇒ Worldwide effort to get high precision data (**ELSA**, JLab, MAMI,...)



# Why study $\pi^0$ and $\eta$ in the final state?

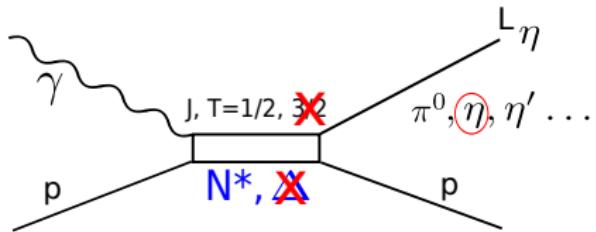
## $\pi^0$ -photoproduction

- high cross section  
→ Large statistics



## $\eta$ -photoproduction

- $\eta$  ( $T=0$ ) → exclusive access to intermediate states  $N^*$  with  $T=1/2$
- low contributions from non-resonant terms



# Which observables to measure?

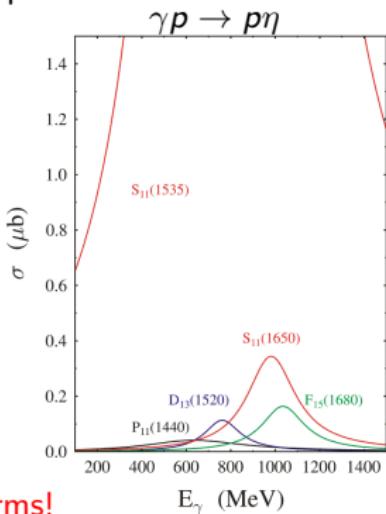
- Scattering amplitude  $f \longleftrightarrow 4$  complex amplitudes (CGLN amplitudes)  
 $f(F_1(W, \cos \theta_{cm}), F_2(W, \cos \theta_{cm}), F_3(W, \cos \theta_{cm}), F_4(W, \cos \theta_{cm}))$
- PWA:  $F_1 = \sum_{I=0}^{\infty} (IM_{I+} + E_{I+})P'_{I+1} + [(I+1)M_{I-} + E_{I-}]P'_{I-1}$ 
  - $E_{I\pm}(W), M_{I\pm}(W)$ : Multipoles
  - $P'_{I\pm 1}(\cos \theta_{cm})$ : Legendre polynomials
- Measurable observables  $\longleftrightarrow$  Multipoles  $\longleftrightarrow$  Resonance parameters

Photon polarization	Target polarization			Recoil nucleon polarization				Target and recoil polarizations			
	X	Y	Z <sub>(beam)</sub>	X'	Y'	Z'	X'	X'	Z'	Z'	
				X	Z	X	Z	X	Z	X	Z
unpolarized	$\sigma$	-	T	-	-	P	-	$T_x$	$L_x$	$T_z$	$L_z$
linear	$-\Sigma$	H	(-P)	-G	$O_x$	(-T)	$O_z$	(-L <sub>x</sub> )	(T <sub>z</sub> )	(L <sub>x</sub> )	(-T <sub>x</sub> )
circular	-	F	-	-E	$C_x$	-	$C_z$	-	-	-	-

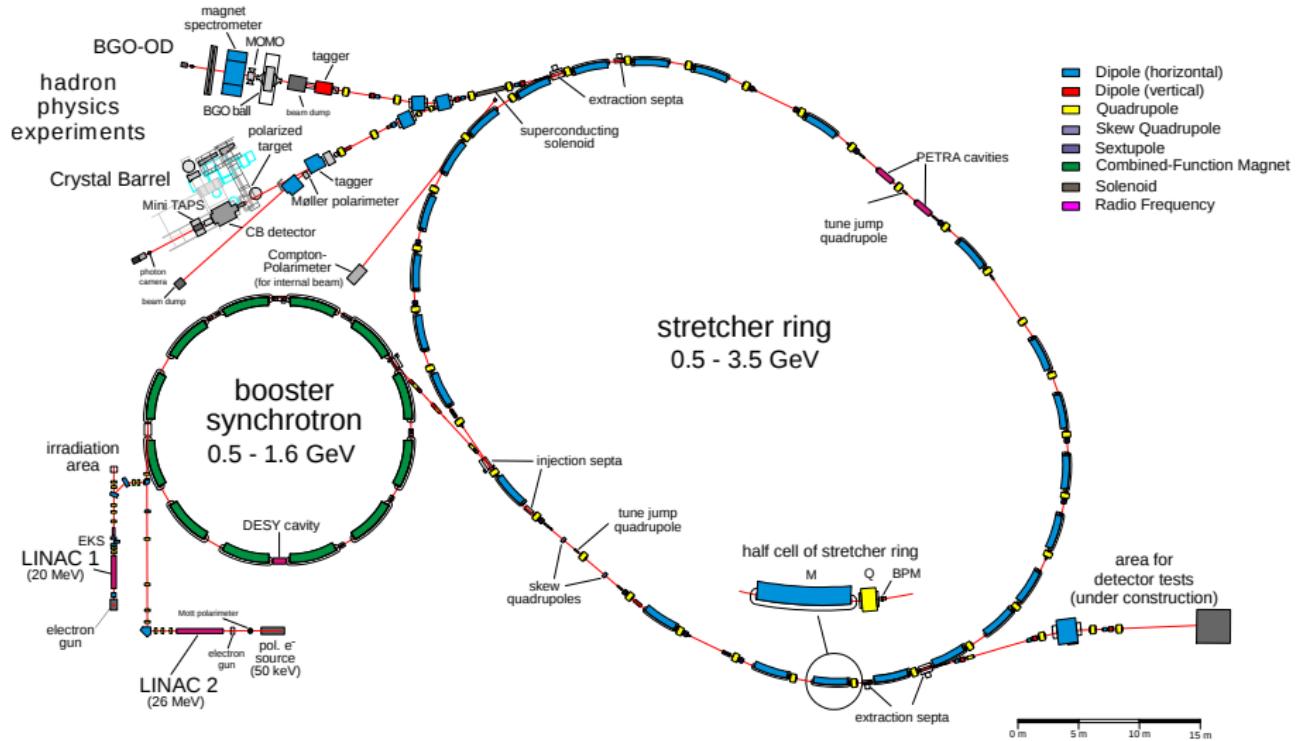
$$\sigma \sim |E_{0+}|^2 + |E_{1+}|^2 + |M_{1+}|^2 + |M_{1-}|^2 + \dots$$

$$\Sigma \sim -2E_{1+}^* M_{1+} + 2M_{1-}^* E_{1+} - 2M_{1-}^* M_{1+} + \dots$$

⇒ Polarization observables are sensitive to interference terms!



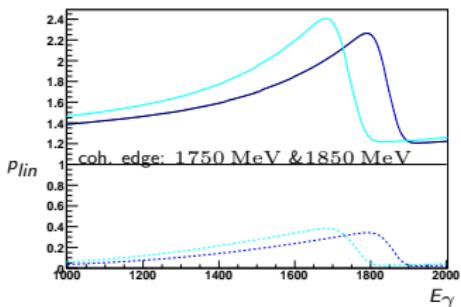
# The Electron Stretcher Accelerator (ELSA)



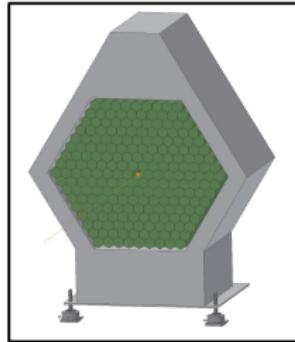
# The CBELSA/TAPS experiment at ELSA in Bonn

Measurement of  $\Sigma$  (July-October 2013)

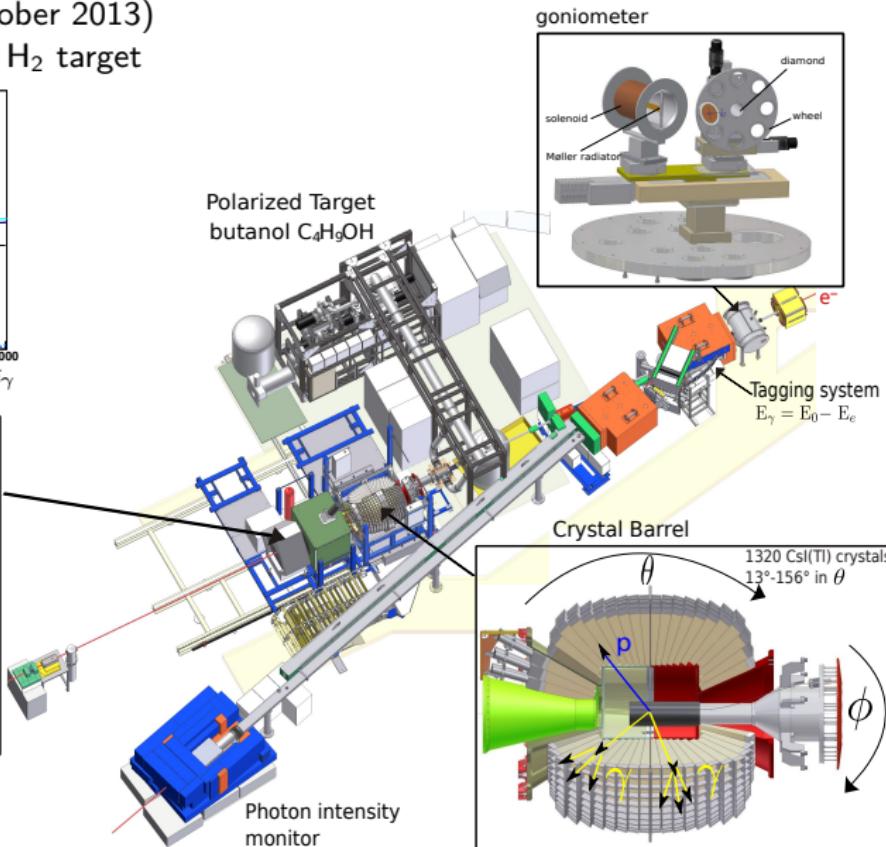
Linearly polarized photons + H<sub>2</sub> target



MiniTAPS



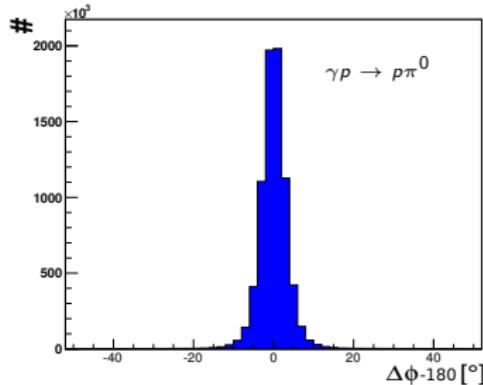
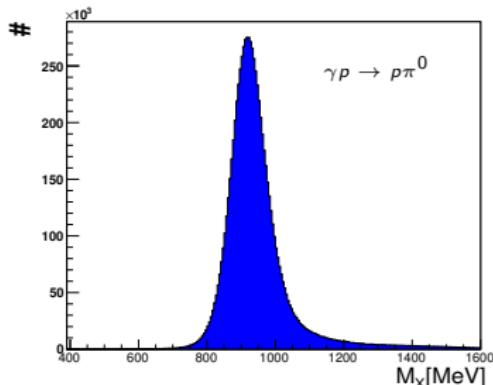
216 BaF<sub>2</sub> crystals  
1°-12° in



# Selection process of $\gamma p \rightarrow \gamma\gamma p$

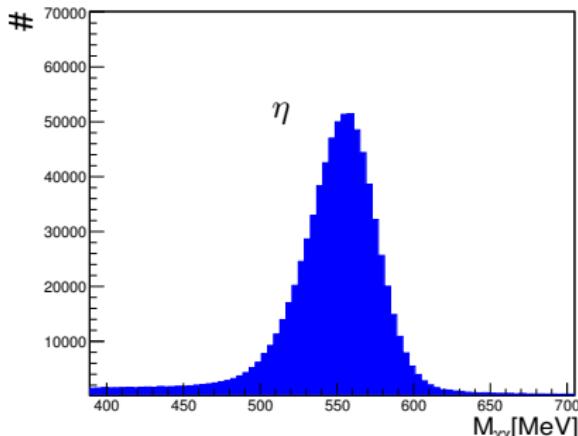
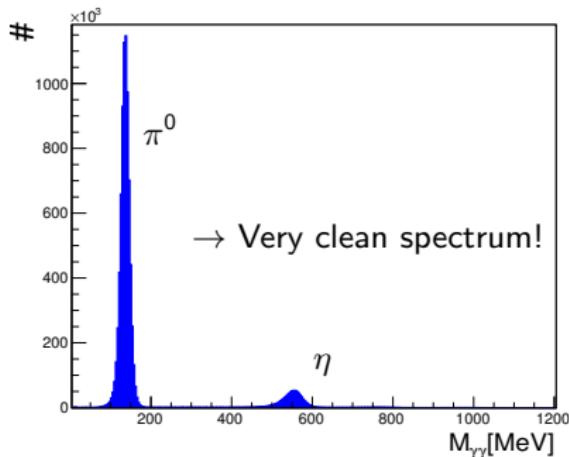
Selected events had to fulfill kinematic constraints:

- 3 hits in calorimeters ( $p+2\gamma$ )
- Proton: calculated as missing particle of  $\gamma p \rightarrow \gamma\gamma X$
- Angular-cuts:
  - Agreement of missing mass and measured charged particle in  $\theta$
  - Coplanarity-cut:  $\Delta\Phi = |\Phi_{\gamma\gamma} - \Phi_p| = 180^\circ$  within  $2.5\sigma$
- Beam photon:  $E_\gamma > E_{\text{prod. threshold}}$  and time coincidence with reaction products



# Selection process of $\gamma p \rightarrow \gamma\gamma p$

- The  $\gamma\gamma$  invariant mass:

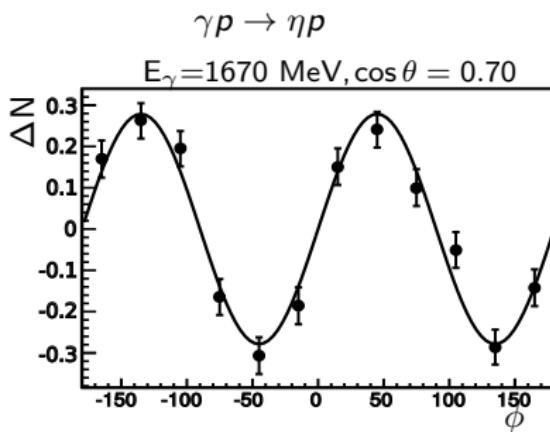
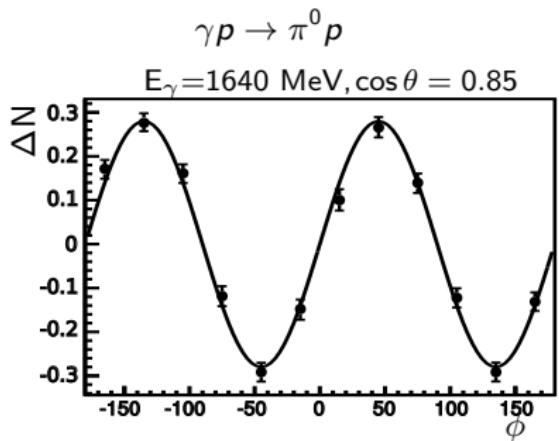
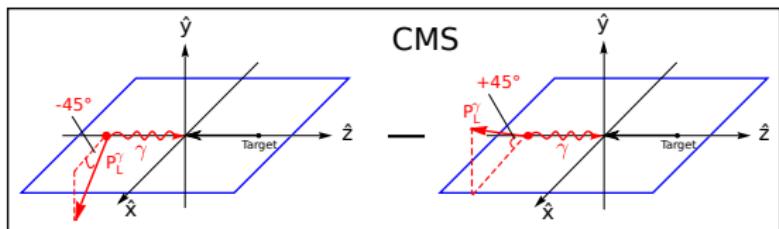


- $5.4 \cdot 10^6 \pi^0$ -events were selected
- $6.6 \cdot 10^5 \eta$ -events were selected

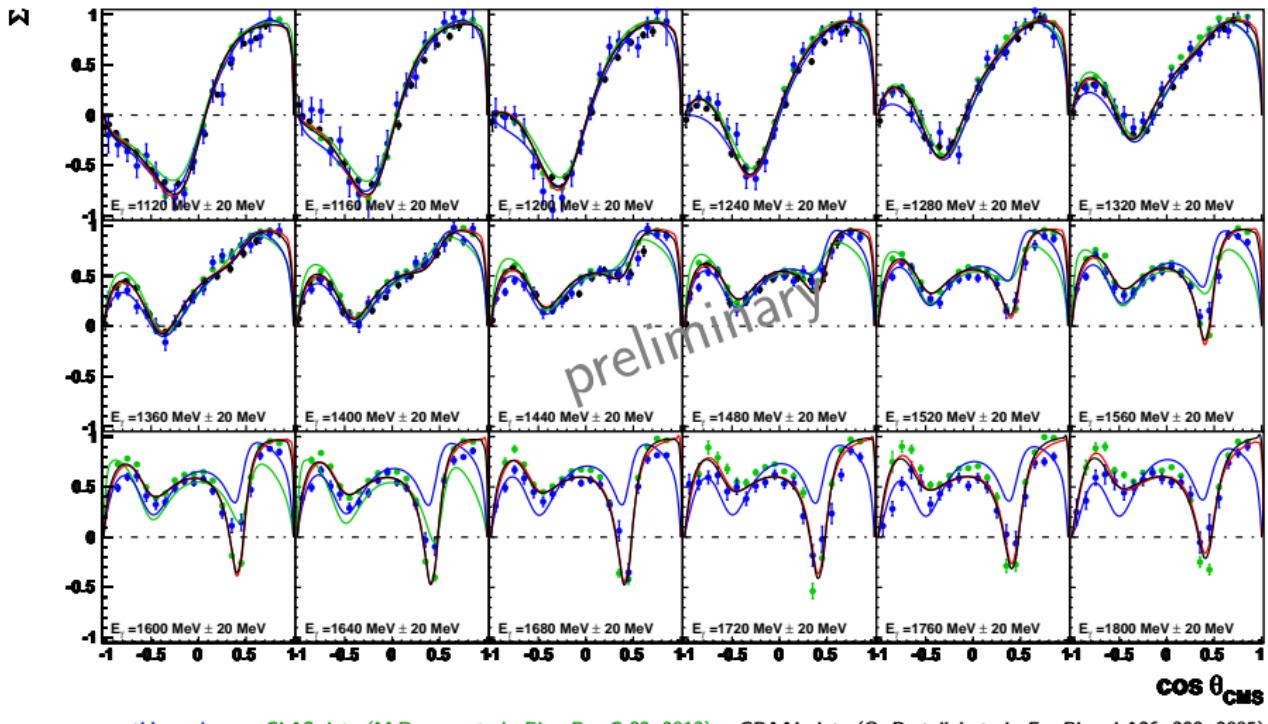
# Determination of the beam asymmetry $\Sigma$

- linearly polarized beam, unpolarized liquid hydrogen target

$$\Delta N = \frac{N_{-45^\circ} - N_{+45^\circ}}{N_{-45^\circ} + N_{+45^\circ}} = p_\gamma^{\text{lin}} \Sigma \sin(2\phi)$$



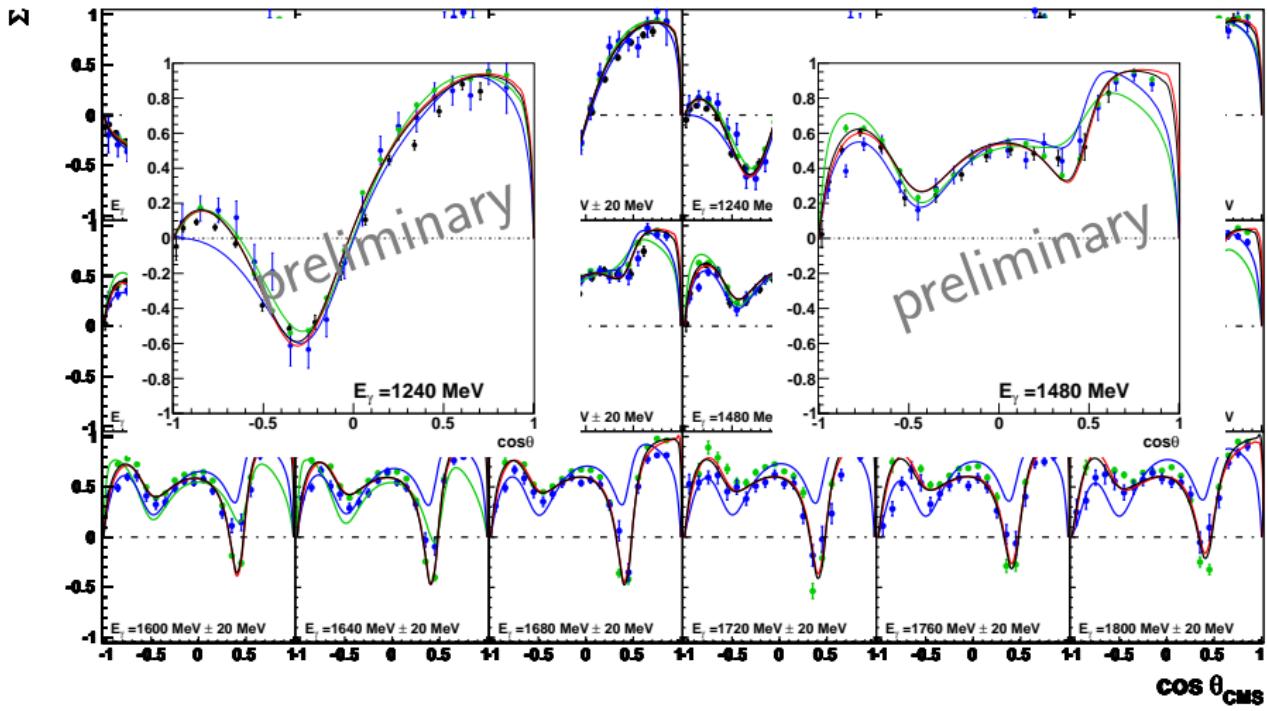
# The beam asymmetry $\Sigma$ in $\pi^0$ -photoproduction



• this work • CLAS data (M.Dugger et al., Phys.Rev.C 88, 2013) • GRAAL data (O. Bartalini et al., Eur.Phys.J.A26, 399, 2005)

PWA solutions: —BnGa(2014\_01) —BnGa(2014\_02) —SAID(CM12) —MAID2007

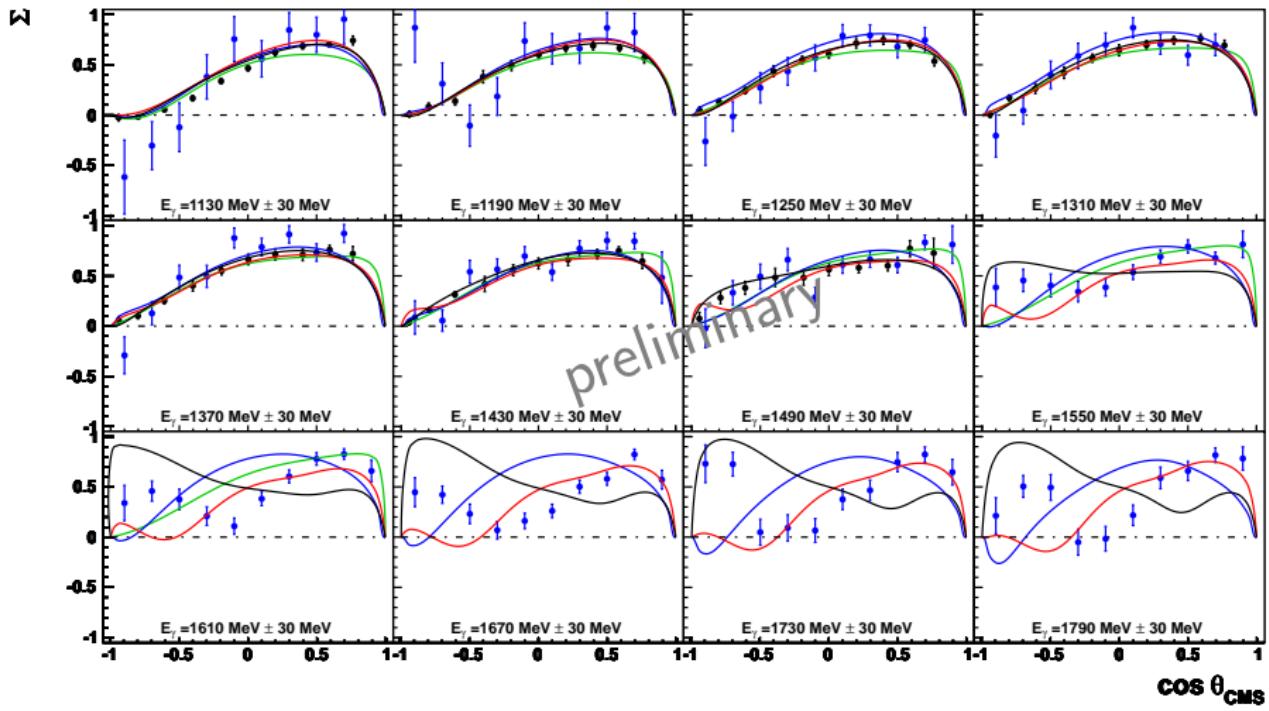
# The beam asymmetry $\Sigma$ in $\pi^0$ -photoproduction



● this work     ● CLAS data (M.Dugger et al., Phys.Rev.C 88, 2013)     ● GRAAL data (O. Bartalini et al., Eur.Phys.J.A26, 399, 2005)

PWA solutions: —BnGa(2014\_01)   —BnGa(2014\_02)   —SAID(CM12)   —MAID2007

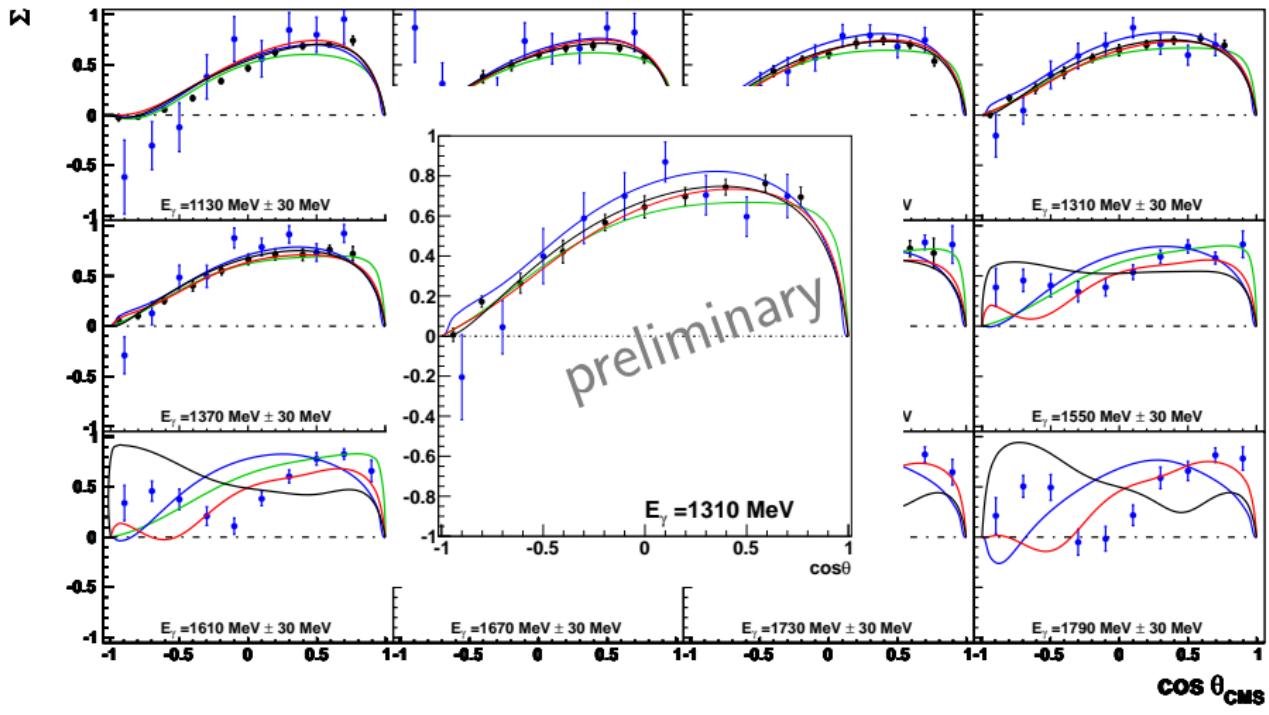
# The beam asymmetry $\Sigma$ in $\eta$ -photoproduction



● this work ● GRAAL data (O.Bartalini et al., Eur.Phys.J. A33, 169, 2007)

PWA solutions: —BnGa(2014\_01) —BnGa(2014\_02) —SAID(GE09) — $\eta$ -MAID

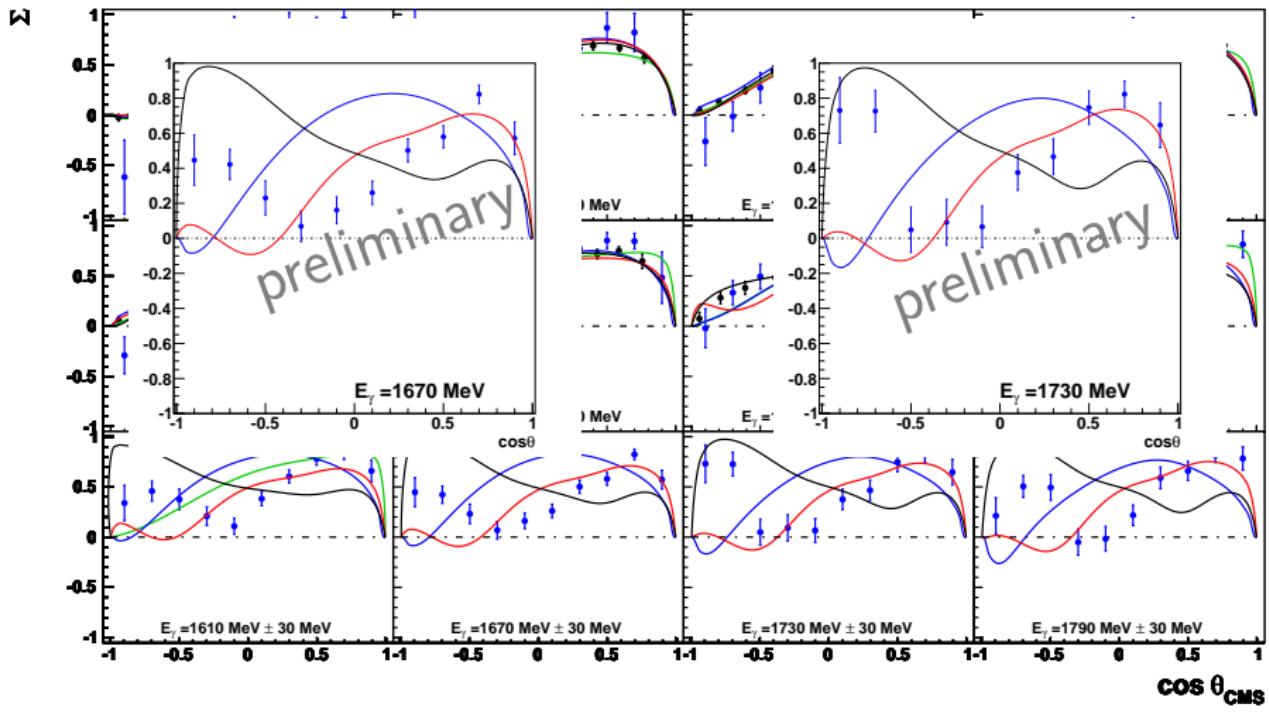
# The beam asymmetry $\Sigma$ in $\eta$ -photoproduction



● this work ● GRAAL data (O.Bartalini et al., Eur.Phys.J. A33, 169, 2007)

PWA solutions: —BnGa(2014\_01) —BnGa(2014\_02) —SAID(GE09) — $\eta$ -MAID

# The beam asymmetry $\Sigma$ in $\eta$ -photoproduction



● this work ● GRAAL data (O.Bartalini et al., Eur.Phys.J. A33, 169, 2007)

PWA solutions: —BnGa(2014\_01) —BnGa(2014\_02) —SAID(GE09) — $\eta$ -MAID

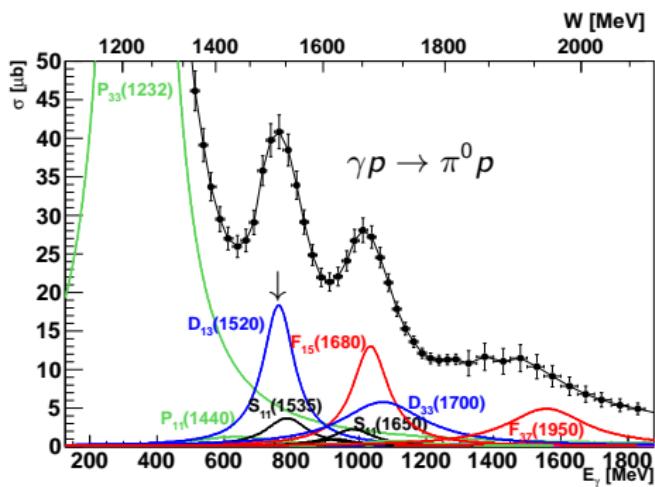
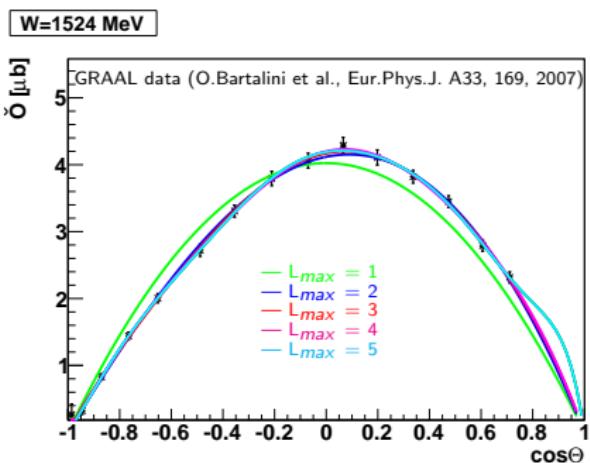
- The beam asymmetry  $\Sigma$  was determined in  $\pi^0$ - and  $\eta$ -photoproduction
- Results:
  - very precise  $\pi^0$  data was measured for  $E_\gamma = 1100 \text{ MeV} - 1800 \text{ MeV}$
  - precise  $\eta$  data was measured for  $E_\gamma = 1100 \text{ MeV} - 1800 \text{ MeV}$
  - $\eta$  data can not be described by either PWA models
  - data will provide new constraints for the PWA
- More results in other polarization observables  
→ See next talk by Jonas Müller

Thank you!

# Truncated PWA (which $L_{max}$ is seen in the data?)

$$\hat{\Sigma}(W, \cos \theta) = \Sigma(W, \cos \theta) \cdot \frac{d\sigma}{d\Omega}(W, \cos \theta) = \sum_{k=2}^{2+2L_{max}} (a_L(W))_k \cdot P_k^2(\cos \theta)$$

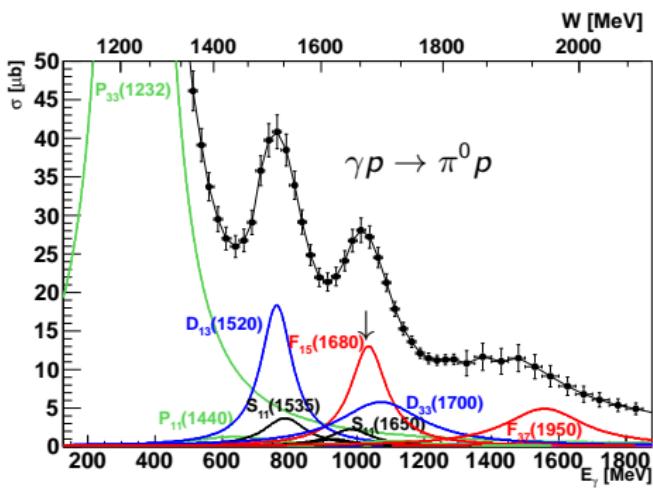
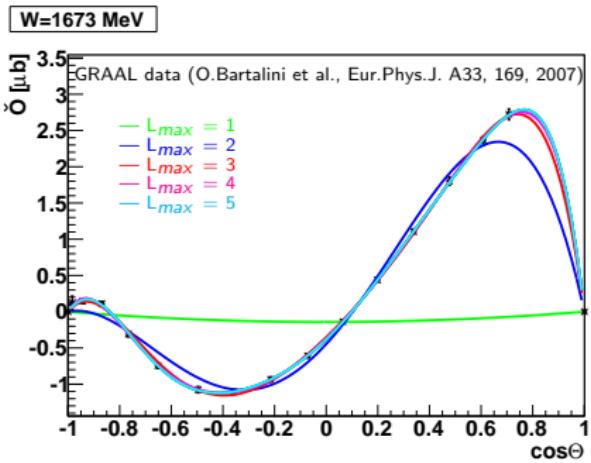
$L_{max} = 0$ S-wave	$L_{max} = 1$ P-wave	$L_{max} = 2$ D-wave	$L_{max} = 3$ F-wave	$L_{max} = 4$ G-wave
$S_{11}(1535)$	$P_{11}(1440)$	$D_{13}(1520)$	$F_{15}(1680)$	$G_{17}(2190)$
$S_{11}(1650)$	$P_{13}(1720)$	$D_{15}(1675)$	$F_{35}(1905)$	
$S_{31}(1620)$	$P_{33}(1232)$	$D_{33}(1700)$	$F_{37}(1950)$	



# Truncated PWA (which $L_{max}$ is seen in the data?)

$$\hat{\Sigma}(W, \cos \theta) = \Sigma(W, \cos \theta) \cdot \frac{d\sigma}{d\Omega}(W, \cos \theta) = \sum_{k=2}^{2+2L_{max}} (a_L(W))_k \cdot P_k^2(\cos \theta)$$

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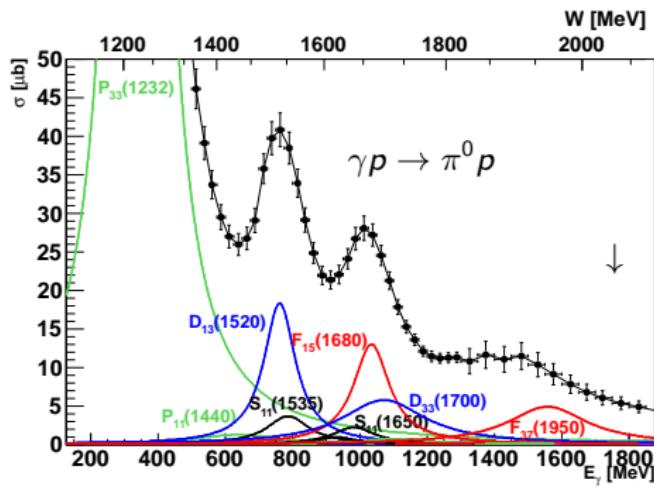
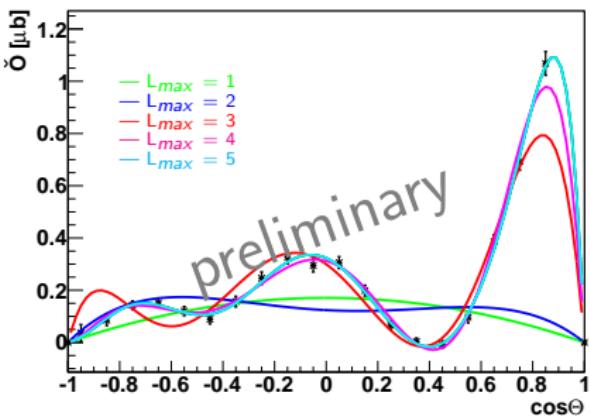


# Truncated PWA (which $L_{max}$ is seen in the data?)

$$\hat{\Sigma}(W, \cos \theta) = \Sigma(W, \cos \theta) \cdot \frac{d\sigma}{d\Omega}(W, \cos \theta) = \sum_{k=2}^{2+2L_{max}} (a_L(W))_k \cdot P_k^2(\cos \theta)$$

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$S_{11}(1650)$	$P_{13}(1720)$	$D_{15}(1675)$	$F_{35}(1905)$	$G_{37}(2200)?$
$S_{31}(1620)$	$P_{33}(1232)$	$D_{33}(1700)$	$F_{37}(1950)$	

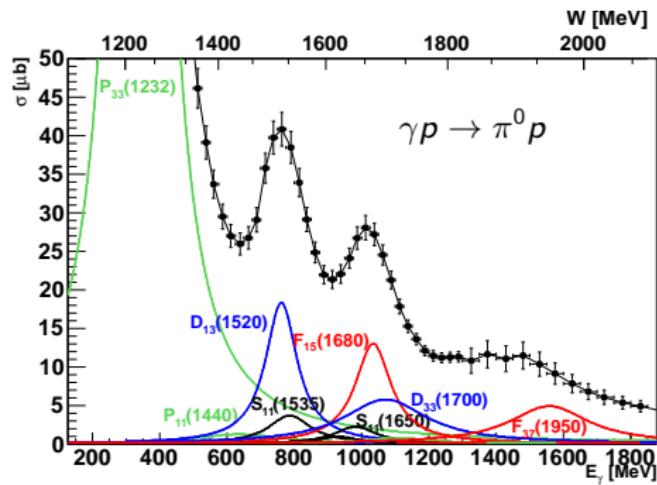
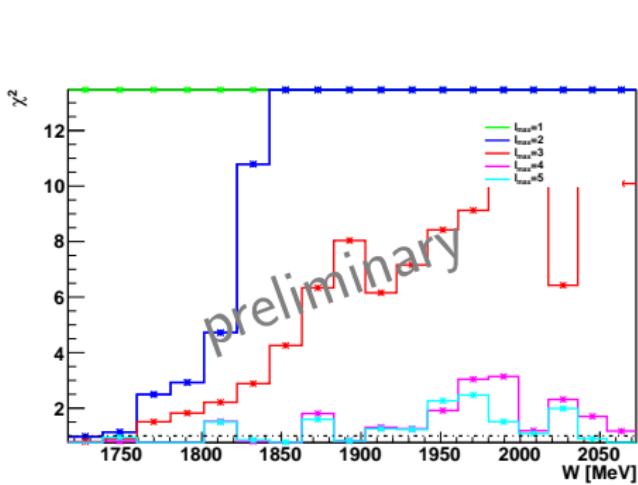
W=2045 MeV



# Truncated PWA (which $L_{max}$ is seen in the data?)

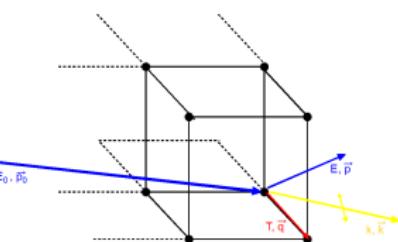
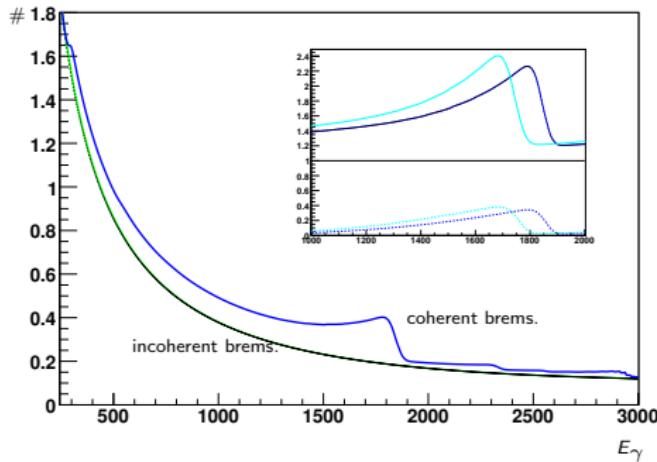
$$\hat{\Sigma}(W, \cos \theta) = \Sigma(W, \cos \theta) \cdot \frac{d\sigma}{d\Omega}(W, \cos \theta) = \sum_{k=2}^{2+2L_{max}} (a_L(W))_k \cdot P_k^2(\cos \theta)$$

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$S_{31}(1620)$	$P_{33}(1232)$	$D_{33}(1700)$	$F_{37}(1950)$	



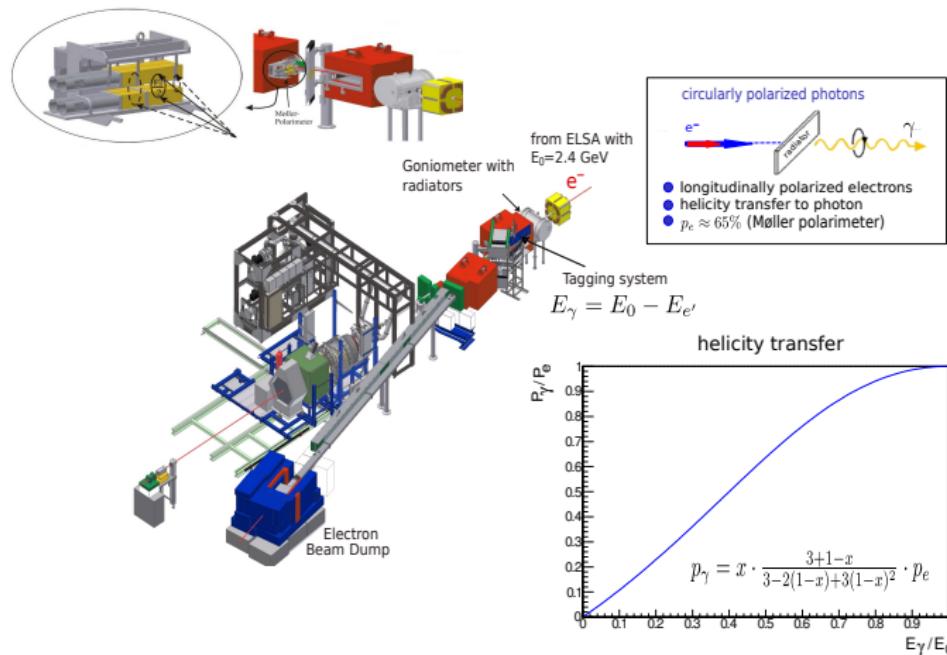
# Linearly polarized photons

- coherent bremsstrahlung produced on diamond crystal
- Bragg: if  $\vec{q} = n \cdot \vec{g}$   $\rightarrow$  constructive interference



# Circularly polarized photons

- Need longitudinally polarized electrons
- helicity transfer from electrons to photons



## Frozen-spin butanol target

- polarize electrons (2.5 T, 300 mK)
- transfer polarization to the protons dynamically via irradiation of microwaves
- "freeze spin" (70 mK) → long relaxation times

